

See discussions, stats, and author profiles for this publication at: <https://www.researchgate.net/publication/11808957>

# The life span of the biosphere revisited

Article in *Nature* · December 1992

DOI: 10.1038/360721a0 · Source: PubMed

CITATIONS

154

READS

276

2 authors:



**Ken Caldeira**

Carnegie Institution for Science

382 PUBLICATIONS 21,147 CITATIONS

[SEE PROFILE](#)



**James F Kasting**

Pennsylvania State University

296 PUBLICATIONS 19,785 CITATIONS

[SEE PROFILE](#)

Some of the authors of this publication are also working on these related projects:



Nitrogen Deposition and Carbon uptake [View project](#)



Carbon Cyle [View project](#)

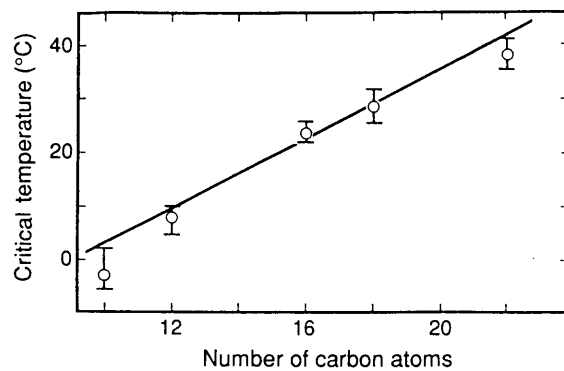


FIG. 4 Threshold temperature  $T_c$  for optimum grafting against the chain length  $n$  of the  $n$ -alkyltrichlorosilane. A linear variation of  $T_c$  is observed with a slope of  $\sim 3.5$  °C per additional methylene group.

above or below a characteristic temperature  $T_0$ , the 'triple point'<sup>13</sup>. This temperature varies linearly with the alkyl chain length<sup>14</sup>, and it is striking that values published for fatty acids fall in the region as in the present experiments. If such an interpretation turns out to be correct, it would suggest that a lateral rearrangement occurs within the monolayer, following physisorption onto the substrate, and preceding chemical grafting. □

Received 18 June; accepted 10 November 1992.

- Bigelow, W. C., Pickett, D. L. & Zisman, W. A. *J. Colloid Sci.* **1**, 513 (1946).
- Plueddeman, E. P., in *Silane Coupling Agents* (Plenum, New York, 1982).
- Maoz, R. & Sagiv, J. *J. Colloid. Interf. Sci.* **100**, 465 (1984).
- Netzer, L., Iscovic, R. & Sagiv, J. *Thin Solid Films* **100**, 67 (1983).
- Gun, J. & Sagiv, J. *J. Colloid. Interf. Sci.* **112**, 457 (1986).
- Tidswell, I. M. *et al. J. chem. Phys.* **95**, 2854 (1991).
- Tidswell, I. M. *et al. phys. Rev.* **B41**, 1111 (1990).
- Wasserman, S. R., Tao, Y. T. & Whitesides, G. M. *Langmuir* **5**, 1074 (1989).
- Silberzan, P., Léger, L., Aussère, D. & Bennattar, J. *J. Langmuir* **7**(8), 1647 (1991).
- Vig, J. R. *J. Vac. Sci. Technol.* **A3**, 1027 (1985).
- Kruger, A. A. in *Surface and Near Surface Chemistry of Oxide Materials* (eds Nowotny, J. & Dufour, L. C.) 413–448 (Elsevier, Amsterdam, 1985).
- Zisman, W. A. *Adv. Chem. Ser. No.* **43**, 1 (1964).
- Gaines, G. L. in *Insoluble Monolayers at Liquid-Gas Interfaces* (Wiley, New York, 1966).
- Kellner, B. M. J., Muller-Landau, F. & Cadenhead, D. A. *J. Colloid. Interf. Sci.* **66**, 3 (1978).
- Allain, C., Aussère, D. & Rondelez, F. *J. Colloid. Interf. Sci.* **107**, 5 (1985).
- Dussan, V. E. *J. Fluid Mech.* **151**, 1 (1985).
- Brochard-Wyart, F., Hervet, H., Redon, C. & Rondelez, F. *J. Colloid. Interf. Sci.* **142**, 518 (1991).

ACKNOWLEDGEMENTS. We thank F. Brochard-Wyart, D. Allara, D. Chatenay and L. Bourdieu for discussions. This work was supported by the Direction des Recherches et Etudes techniques and by the CNRS.

## The life span of the biosphere revisited

Ken Caldeira & James F. Kasting

Earth System Science Center & Department of Geosciences,  
The Pennsylvania State University, University Park,  
Pennsylvania 16802, USA

A DECADE ago, Lovelock and Whitfield<sup>1</sup> raised the question of how much longer the biosphere can survive on Earth. They pointed out that, despite the current fossil-fuel induced increase in the atmospheric  $\text{CO}_2$  concentration, the long-term trend should be in the opposite direction: as increased solar luminosity warms the Earth, silicate rocks should weather more readily, causing atmospheric  $\text{CO}_2$  to decrease. In their model<sup>1</sup>, atmospheric  $\text{CO}_2$  falls below the critical level for C3 photosynthesis, 150 parts per million (p.p.m.), in only 100 Myr, and this is assumed to mark the demise of the biosphere as a whole. Here, we re-examine this problem using a more elaborate model that includes a more accurate treatment of the greenhouse effect of  $\text{CO}_2$  (refs 2–4), a biologically mediated weathering parameterization, and the realization that

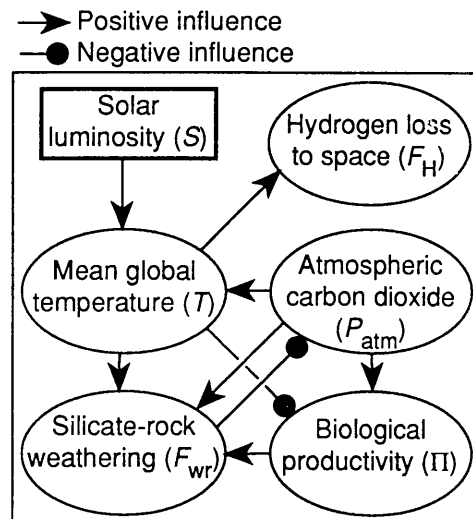


FIG. 1 Diagram illustrating the positive and negative influences represented in our model. System behaviour may be understood by examining the response to an increase in solar luminosity ( $S$ ). Increasing  $S$  warms the Earth, increasing  $T$ . This enhances both the weatherability of silicate rocks ( $F_{wr}$ ) and the rate of hydrogen escape to space ( $F_H$ ). At temperatures approaching 50 °C, temperature increases have a negative influence on biological productivity ( $\Pi$ ). Enhanced silicate-rock weatherability draws down the atmospheric  $\text{CO}_2$  concentration ( $P_{atm}$ ). The lower  $\text{CO}_2$  concentration tends to buffer  $T$  and  $F_{wr}$  and to reduce  $\Pi$ . Eventually, either lack of  $\text{CO}_2$ , high temperatures, or the loss of water will limit the life span of the biosphere.

C4 photosynthesis can persist to much lower concentrations of atmospheric  $\text{CO}_2$  (<10 p.p.m.)<sup>5,6</sup>. We find that a C4-plant-based biosphere could survive for at least another 0.9 Gyr to 1.5 Gyr after the present time, depending respectively on whether  $\text{CO}_2$  or temperature is the limiting factor. Within an additional 1 Gyr, Earth may lose its water to space, thereby following the path of its sister planet, Venus.

The problem of the life span of the biosphere has implications not only for the future of our planet, but also for the probability of finding biologically active planets in our galactic neighbourhood. As the Sun converts hydrogen to helium, its core becomes denser and hotter, increasing the rate of thermonuclear fusion; hence, standard solar models predict that the Sun has been getting more luminous with time<sup>7–9</sup>. As solar luminosity increases, silicate rocks should weather more easily, thereby drawing down  $\text{CO}_2$  from the Earth's atmosphere. This feedback mechanism should tend to buffer the Earth's temperature near its present value, both in the future and in the past<sup>10</sup>. Eventually, the  $\text{CO}_2$  concentration may become so low that the megafloa existing at present will not be able to engage profitably in photosynthesis, effectively cutting off the carbon supply for the biosphere<sup>1</sup>. Vanishingly low  $\text{CO}_2$  concentrations would preclude further  $\text{CO}_2$ -modulated thermal buffering. The Earth would then warm more rapidly, and much of the remaining biota might be pressed against thermal barriers to their survival. Ultimately, as the Sun continues to grow brighter, the Earth's surface water will be lost through photodissociation and escape of hydrogen to space. The loss of surface water would bring an indisputable end to the lifetime of the biosphere.

We developed a computer model (see Box 1 and Fig. 1) to analyse the situation described above. Because our model maximizes both the greenhouse effect and stratospheric water loss, and because we have estimated temperature and carbon requirements for the biosphere conservatively, our model should produce a minimum estimate for the biosphere's life span. We assume that future  $\text{CO}_2$  emissions from volcanic and mid-ocean-ridge sources will continue at roughly the present rate, although there is some possibility that these rates will increase in the future<sup>11</sup>. Thus, as long as most of the calcium and magnesium

BOX 1 Model description

Our computer model couples solar luminosity ( $S$ ), the silicate-rock weathering rate ( $F_{wr}$ ), and the global energy balance, to estimate the partial pressure of atmospheric carbon dioxide ( $P_{atm}$ ), and biological productivity ( $\Pi$ ) as a function of time ( $t$ ) into the future.

The luminosity ( $S$ ) of the Sun on the main sequence as measured at the Earth was calculated for this work by A. I. Boothroyd (personal communication) using a standard solar model<sup>9</sup> with updated physical constants and  ${}^7\text{Be}(e^-, \nu){}^7\text{Li}$  reaction rate. A curve fit to these model results yields

$$S(t) = (1 - 0.38t/\tau_0)^{-1} S_0 \quad (1)$$

with a mean error <0.6% for the interval  $-4.5 \text{ Gyr} < t < 4.77 \text{ Gyr}$ . Here  $t$  is time expressed as years from the present,  $\tau_0 = 4.55 \text{ Gyr}$ , and the subscript 0 refers to present-day values ( $S_0 = 1.368 \text{ W m}^{-2}$ ). An energy balance may be written for the Earth by setting incoming shortwave radiation equal to the outgoing longwave radiation

$$(1 - a)S/4 = \sigma T_{eff}^4 \quad (2)$$

where  $a$  is the planetary albedo,  $\sigma$  is the Stefan-Boltzmann constant, and  $T_{eff}$  is the effective black-body radiation temperature of the Earth.  $T$  may be related to  $T_{eff}$  by the expression

$$T = T_{eff} + \Delta T \quad (3)$$

The greenhouse warming factor ( $\Delta T$ ) is a function of  $T$  (in K) and  $\psi = \log P_{atm}$  (in bars):

$$\Delta T = 815.17 + (4.895 \times 10^7)T^{-2} - (3.9787 \times 10^6)T^{-1} - 6.7084\psi^{-2} + 73.221\psi^{-1} - 30.882T^{-1}\psi^{-1} \quad (4)$$

The planetary albedo ( $a$ ) is a function of  $T$

$$a = 1.4891 - 0.0065979T + (8.567 \times 10^{-6})T^2 \quad (5)$$

Equations (4) and (5) were developed using least-square fits to the results of 143 runs of a radiative-convective climate model<sup>2-4</sup> with  $0^\circ\text{C} \leq T \leq 100^\circ\text{C}$  and  $10^{-8} \text{ bar} \leq P_{atm} \leq 10^{-2} \text{ bar}$ ; the mean errors were <0.5 K and <0.01 for equations (4) and (5), respectively. This radiative-convective climate model<sup>2-4</sup> has an  $\text{H}_2\text{O}$ -saturated troposphere, producing maximum infrared trapping and maximum stratospheric water loss. Hence, this model should generate lower bounds on the potential life span of the biosphere. Because it is not known how clouds will respond to surface temperature changes, the effect of clouds on shortwave radiation forcing ( $14 \text{ W m}^{-2}$ )

and planetary albedo were held constant. We consider this to be a conservative assumption because, for large perturbations, clouds would probably act to stabilize Earth's temperature<sup>16</sup>.

The activity of  $\text{H}^+$  in fresh soil water ( $a_{H^+}$ ) can be calculated assuming equilibration of rain water with the soil  $\text{CO}_2$  concentration ( $P_{soil}$ ) and the atmospheric  $\text{SO}_2$  concentration (0.2 p.p.b.). Equilibrium constants and relations for the carbon and sulphur systems may be found in refs 14 and 23, respectively. Experimental results in the acid pH range at  $25^\circ\text{C}$  with low  $\text{CO}_2$  concentrations indicate that silicate dissolution rates are typically proportional to the hydrogen-ion activity to roughly the 0.5 power<sup>24,25</sup>. Hence, incorporating the temperature dependency proposed by Walker *et al.*<sup>10</sup>, the silicate-rock weathering rate ( $F_{wr}$ ), may be represented

$$\frac{F_{wr}}{F_{wr,0}} = \left( \frac{a_{H^+}}{a_{H^+,0}} \right)^{0.5} \exp \left( \frac{T - T_0}{13.7 \text{ K}} \right) \quad (6)$$

The exponent in equation (6) is in good agreement with the original formulation of Walker *et al.*<sup>10</sup> (with an implied exponent of 0.6), developed using experimental results<sup>26,27</sup> obtained at 100 to  $200^\circ\text{C}$  and 2 to 20 bars  $\text{CO}_2$ .

Volk<sup>28</sup> parameterized  $P_{soil}$  as a function of  $\Pi$  and  $P_{atm}$

$$\frac{P_{soil}}{P_{soil,0}} = \frac{\Pi}{\Pi_0} \left( 1 - \frac{P_{atm,0}}{P_{soil,0}} \right) + \frac{P_{atm}}{P_{soil,0}} \quad (7)$$

We extended Volk's<sup>28</sup> representation of  $\Pi$  by including a multiplicative, temperature-dependent term that produces maximum productivity when  $T$  is  $25^\circ\text{C}$  and zero productivity when  $T$  reaches  $50^\circ\text{C}$

$$\frac{\Pi}{\Pi_{max}} = \left( 1 - \left( \frac{T - 25^\circ\text{C}}{25^\circ\text{C}} \right)^2 \right) \left( \frac{P_{atm} - P_{min}}{P_{1/2} + (P_{atm} - P_{min})} \right) \quad (8)$$

A maximum possible biological productivity ( $\Pi_{max}$ ) equal to twice the present productivity ( $\Pi_0$ ) was deemed by Volk<sup>28</sup> to be the most defensible value, and is adopted for this study. The minimum  $\text{CO}_2$  concentration under which C4 plants can grow ( $P_{min}$ ) was conservatively estimated to be 10 p.p.m. The value of  $P_{1/2}$  is calculated by forcing equation (8) to yield  $\Pi = \Pi_0$  when  $T = T_0 (=15^\circ\text{C})$  and  $P_{atm} = P_{atm,0} (=320 \text{ p.p.m.})$ . Following Volk<sup>28</sup>, we use  $P_{soil,0} = 10 P_{atm,0}$  as a representative soil  $p\text{CO}_2(P_{soil})$ .

The model is solved by setting  $F_{wr} = F_{wr,0}$  until  $P_{atm} = 1 \text{ p.p.m.}$ ; thereafter,  $P_{atm}$  is maintained at this value.

released in the weathering process accumulates in carbonate sediments, future weathering rates should be roughly the same as today's.

A principal difference between our model and that of Lovelock and Whitfield<sup>1</sup> is that our climate model is much more sensitive to  $\text{CO}_2$  decreases below the present level of 350 p.p.m. Their proposed constant temperature relation between the atmospheric  $\text{CO}_2$  partial pressure ( $P_{atm}$ ) and solar luminosity ( $S$ ) ( $P_{atm}^{0.5} \propto 1.017 S_0 - S$ ) implies that the infrared trapping by 320 p.p.m.  $\text{CO}_2$  is only  $\sim 4 \text{ W m}^{-2}$ . We calculate however, that the actual amount of infrared trapping by this concentration of  $\text{CO}_2$  is  $20.4 \text{ W m}^{-2}$ , using a radiative-convective model<sup>2-4</sup> at  $15^\circ\text{C}$  with relative humidity as specified by Manabe and Wetherald<sup>12</sup>. This is about six times the radiative forcing ( $3.4 \text{ W m}^{-2}$ ) that would result from a  $\text{CO}_2$  doubling in this model. (The absorption of infrared radiation increases roughly logarithmically with the  $\text{CO}_2$  concentration because the strongest absorption bands, for example the  $15\text{-}\mu\text{m}$  band, are already saturated near their centres.) If  $S$  increases with time according to equation (1) ( $\sim 0.12 \text{ W m}^{-2} \text{ Myr}^{-1}$ ), then reductions in atmospheric  $\text{CO}_2$  concentration could maintain a constant temperature for the next 0.9 Gyr. If we use the critical value for photosynthesis suggested by Lovelock and Whitfield ( $P_{min} = 150 \text{ p.p.m.}$ ), our model calculates that 0.5 Gyr of the biosphere's life span remains, five times the Lovelock and Whitfield estimate. But, Lovelock and Whitfield ignored C4 plants, some of which can survive at  $\text{CO}_2$  concentrations below 10 p.p.m. (refs 5 and 6). Using this criterion ( $P_{min} = 10 \text{ p.p.m.}$ ), our model calculates that the remaining life span of the biosphere is 0.9 Gyr (Fig. 2).

There is reason to suspect that  $\text{CO}_2$  starvation may not limit the life span of the biosphere. As solar luminosity continues to

increase, atmospheric  $\text{CO}_2$  should continue to be drawn down; however, it may reach a minimum at which some C4 plants could survive. Equation (6) indicates that when the global mean temperature is above  $300 \text{ K}$ , silicate dissolution may proceed more rapidly than  $\text{CO}_2$  and  $\text{H}_2\text{SO}_4$  are supplied to the atmosphere by volcanism. This would leave an excess of  $\text{Ca}^{2+}$  and  $\text{Mg}^{2+}$  cations which could not be precipitated as carbonates or sulphates. They would probably precipitate in the oceans in a variety of silicate phases. Geological carbon inputs to the ocean and atmosphere must be largely balanced by carbonate sedimentation, suggesting that the oceans will continue to be saturated with respect to some carbonate phase. If future ocean waters are roughly in chemical equilibrium with silicates and carbonates, ocean pH would be buffered close to its present value<sup>13</sup>. If ocean pH increases by no more than half a pH unit and the oceanic  $\text{Ca}^{2+}$  concentration changes by no more than an order of magnitude, then application of carbonate equilibria<sup>14</sup> indicates that atmospheric partial pressure of  $\text{CO}_2$  ( $p\text{CO}_2$ ) would be reduced no more than two orders of magnitude from today's value. Even if ocean chemistry is not this well buffered, Henry's law will partition more  $\text{CO}_2$  into the atmosphere on a warmer Earth, so the atmosphere may still contain a few parts per million of  $\text{CO}_2$ . Some present-day C4 plants could survive at these  $\text{CO}_2$  concentrations<sup>5,6</sup>, and aquatic microbial-based food chains could flourish at even lower  $\text{CO}_2$  concentrations<sup>15</sup>.

If  $\text{CO}_2$  starvation does not limit the life span of the nonmicrobial biosphere, high temperatures may. Even with negligible atmospheric  $\text{CO}_2$ , global mean temperatures in our model exceed  $50^\circ\text{C}$  after  $\sim 1.5 \text{ Gyr}$ . Our model does not include the possible negative feedback of increasing cloud cover<sup>16</sup>; hence 1.5 Gyr should be considered a minimum time at which this

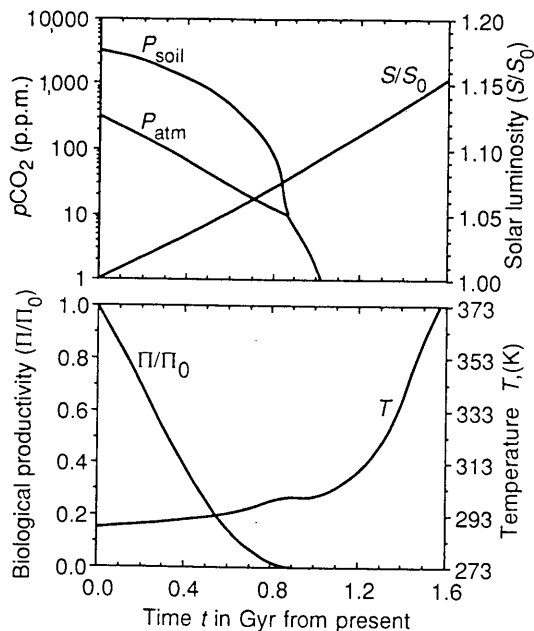


FIG. 2 Model results for the next 1.6 Gyr. The terrestrial biosphere, whose productivity ( $\Pi/\Pi_0$ ) depends on an adequate supply of atmospheric  $\text{CO}_2$ , maintains a gradient between soil and atmospheric  $\text{CO}_2$  concentrations. For  $\sim 1$  Gyr, the partial pressure of soil  $\text{CO}_2$  ( $P_{\text{soil}}$ ) and Earth's mean temperature ( $T$ ) covary such that silicate-rock weathering consumes  $\text{CO}_2$  at the modern rate.  $T$  is a function of both solar luminosity ( $S/S_0$ ) and the atmospheric  $\text{CO}_2$  concentration ( $P_{\text{atm}}$ ). As  $S/S_0$  increases,  $P_{\text{atm}}$  decreases, diminishing  $\Pi/\Pi_0$  and with it the difference between  $P_{\text{soil}}$  and  $P_{\text{atm}}$ . When  $P_{\text{atm}}$  is too low for C4 plants to survive ( $\sim 0.9$  Gyr), then  $P_{\text{soil}} = P_{\text{atm}}$ . (In a C3-only world with  $P_{\text{min}} = 150$  p.p.m., C3 plants survive  $\sim 0.5$  Gyr.) After  $\sim 1$  Gyr, silicate-rock dissolution exceeds volcanic  $\text{CO}_2$  degassing;  $P_{\text{atm}}$  is then held at 1 p.p.m. The timescale for the loss of the oceans through hydrogen escape becomes limited by solar extreme ultraviolet when  $T$  climbs above  $80^\circ\text{C}$  ( $\sim 1.5$  Gyr). The loss of surface water by this process would take  $\sim 1$  Gyr, and could ultimately limit the life span of the biosphere to  $\sim 2.5$  Gyr.

temperature would be reached. Only prokaryotes and protozoa can live much above this temperature<sup>17</sup>. Beyond this time, planetary warming will probably accelerate. Atmospheric water vapour content increases exponentially with surface temperature if the relative humidity remains constant. This water vapour should trap outgoing longwave radiation and decrease the planetary albedo by absorbing solar near-infrared radiation. Both processes should contribute to planetary warming. In our model, the Earth's surface temperature increases from  $50^\circ\text{C}$  to  $100^\circ\text{C}$  in  $<0.2$  Gyr. Some archaeobacteria might survive even higher temperatures<sup>18,19</sup>, but all higher forms of life would certainly be eliminated by this stage.

The final sterilization of the Earth will occur when the planet loses its water. Today, the timescale for the loss of the oceans through photodissociation and hydrogen escape is much longer than the age of the Earth; as the atmosphere warms beyond  $60$  to  $70^\circ\text{C}$ , however, the  $\text{H}_2\text{O}$  mixing ratio in the stratosphere increases markedly<sup>20</sup>. As surface temperatures approach  $\sim 80^\circ\text{C}$ , the stratospheric  $\text{H}_2\text{O}$  mixing ratio reaches  $\sim 2.5\%$ . Above this mixing ratio, the loss of hydrogen to space is limited by the solar extreme ultraviolet heating rate to  $\sim 6 \times 10^{11}$  atoms  $\text{H cm}^{-2} \text{s}^{-1}$  (ref. 21), giving a timescale for ocean loss of  $\sim 1$  Gyr. Hence, the complete elimination of water  $\sim 2.5$  Gyr from now could bring terrestrial life to its final close. After Earth's water is lost, silicate-rock weathering will cease; hence, volcanic  $\text{CO}_2$  should accumulate in the atmosphere, creating a climate much like that of Venus.

Events such as a collision with the giant comet Chiron<sup>22</sup> or a radical reorganization of the carbonate-silicate cycle<sup>11</sup> could possibly bring the life of the biosphere to an earlier close. But given the otherwise conservative assumptions that we have made

in constructing our model, our results indicate that the biosphere will probably survive at least another 1 Gyr, and possibly much longer. This prediction ought to be encouraging to humans, as well, because it implies that we could survive for geologically long time periods if we can manage to cope with our other societal problems. □

Received 31 August; accepted 12 November 1992.

1. Lovelock, J. E. & Whitfield, M. *Nature* **296**, 561-563 (1982).
2. Kasting, J. F. & Ackerman, T. P. *Science* **234**, 1383-1385 (1986).
3. Kasting, J. F. *Paleogeogr. Paleoclimat. Paleocool.* **75**, 83-95 (1989).
4. Kasting, J. F., Whitfield, D. P., & Reynolds, R. T. *Icarus* (in the press).
5. Heath, O. V. S. *The Physiological Aspects of Photosynthesis* (Stanford Univ. Press, (1969).
6. Pearcy, R. W. & Ehleringer, J. *Plant Cell Envir.* **7**, 1-13 (1984).
7. Newman, M. J. & Rood, R. T. *Science* **198**, 1035-1037 (1977).
8. Gough, D. O. *Solar Phys.* **74**, 21-34 (1981).
9. Sackman, I.-J., Boothroyd, A. I. & Fowler, W. A. *Astrophys. J.* **360**, 727-736 (1990).
10. Walker, J. C. G., Hays, P. B. & Kasting, J. F. *J. geophys. Res.* **86**, 9776-9782 (1981).
11. Caldeira, K. *Geology* **19**, 204-206 (1991).
12. Manabe, S. & Wetherald, R. T. *J. Atmos. Sci.* **24**, 241-259 (1967).
13. Sillen, L. G. in *Oceanography* (ed. Sears, M.) 549-581 (Am. Assoc. Adv. Sci., Washington DC, 1961).
14. Stumm, W. & Morgan, J. J. *Aquatic Chemistry* (Wiley, New York, 1981).
15. Miller, A. G., Turpin, D. H. & Carvin, D. T. *Plant Physiol.* **75**, 1064-1070 (1984).
16. Rossow, W. B., Henderson-Sellers, A. & Weinreich, S. K. *Science* **217**, 1245-1247 (1982).
17. Brock, T. D. *Science* **230**, 132-138 (1985).
18. Baross, J. A. & Deming, J. W. *Nature* **303**, 423-426 (1983).
19. Stetter, K. O. in *Thermophiles: General, Molecular, and Applied Microbiology* (ed. Brock, T. D.) 39-74 (Wiley, New York, 1986).
20. Kasting, J. F. *Icarus* **74**, 472-494 (1988).
21. Watson, A. J., Donahue, T. M. & Walker, J. C. G. *Icarus* **48**, 150-166 (1981).
22. Sleep, N. H., Zahnle, K. J., Kasting, J. F. & Morowitz, H. J. *Nature* **342**, 139-142 (1989).
23. Chameides, W. L. *J. geophys. Res.* **89**, 4739-4755 (1984).
24. Blum, A. & Lasaga, A. C. *Nature* **331**, 431-433 (1988).
25. Wogelius, R. A. & Walther, J. V. *Geochim. cosmochim. Acta* **55**, 943-954 (1991).
26. Lagache, M. *Bull. Soc. Franc. Miner. Crist.* **88**, 223-253 (1965).
27. Lagache, M. *Geochim. cosmochim. Acta* **40**, 157-161 (1976).
28. Volk, T. *Am. J. Sci.* **287**, 763-779 (1987).

ACKNOWLEDGEMENTS. We thank A. I. Boothroyd for solar model calculations. K. C. was supported by the NSF under a grant awarded in 1991. J.F.K. was supported by the NASA Exobiology Program.

## Electrical conductivity of carbon-bearing granulite at raised temperatures and pressures

Paul W. J. Glover\* & F. J. Vine†

\* Department of Geological Sciences, University College London, Gower Street, London WC1E 6BT, UK

† School of Environmental Sciences, University of East Anglia, Norwich, Norfolk NR4 7TJ, UK

It has long been recognized that the electrical conductivity of the lower continental crust is anomalously high. Both pore-saturating brines<sup>1-5</sup> and conducting films of carbon at grain boundaries<sup>6-10</sup> have been proposed to explain this, but the evidence remains inconclusive. Here we report measurements of electrical conductivity at high temperatures and pressures<sup>11-13</sup> on samples of carbon-bearing and carbon-free granulites with a range of electrolyte saturations. The application of pressure to nominally dry carbon-free samples reduces the electrical conductivity as a result of a progressive reduction in pore connectivity, whereas the carbon-bearing samples show an increase in conductivity under the same conditions—an effect that we ascribe to reconnection of carbon conduction pathways during compaction. Moreover, we find a greater increase in conductivity with temperature for the carbon-bearing samples. In the light of work indicating that the abundance of carbon in high-grade rocks has been underestimated in the past<sup>7,8</sup>, our results provide strong evidence for the role of carbon in lower-crustal conductivity.

Very few laboratory experiments have been done under fully simulated lower-crustal conditions to estimate the contribution of electrolytes to the *in situ* rock conductivity, and only one study, to our knowledge, examined carbon-bearing rock<sup>14</sup>. Here we examine carbon-bearing and carbon-free granulites under a

# When climate and life finally devolve

It is traditional, as the year draws to a close, to try to predict what the future holds in store. It is less common to ask, as Caldeira and Kasting do in this issue, "how much future?"

NEW calculations by Caldeira and Kasting on page 721 of this issue<sup>1</sup> are about the destiny and ultimate demise of the biosphere — which is both dependent on and vulnerable to the evolution of our main-sequence star, the Sun. Their calculations provide the best estimate to date as to when the Earth's staple of abiotic and biotic processes that draw down atmospheric CO<sub>2</sub> will be unable to offset a rising solar flux.

Predictions of the heat death of the biosphere have traditionally come from astronomers. But ten years ago, Earth scientists Lovelock and Whitfield showed that the biota would face a crisis<sup>2</sup> long before our planet is scorched as the Sun expands into a red-giant phase, in about 5 billion years. The crisis will be a scarcity of atmospheric CO<sub>2</sub>.

The trend over geological time has been for atmospheric levels of CO<sub>2</sub> to decrease. (Current rising levels of CO<sub>2</sub> from industrial and agricultural sources are, by comparison, just a blip on the geological landscape.) Driving this long-term steady decrease has been a combination of factors<sup>3</sup>: the growth of continents, a declining geothermal heat flux, a sequence of evolutionary developments that increase the weatherability of the terrestrial surface, and an increasingly radiant Sun, all of which alter the dynamics of the geochemical system that balances the CO<sub>2</sub> supplied by volcanism with its removal by global weathering. The ultimate consequences of this downward trend were spelled out by Lovelock and Whitfield: in just 100 million years, CO<sub>2</sub> will drop below the minimum level necessary to support photosynthesis.

Caldeira and Kasting have recalculated the life span of the biosphere, using a more refined set of biological and physical inputs. These include a precise greenhouse function, increased photosynthetic efficiencies from C4 plants at reduced CO<sub>2</sub> and allowance for a functional dependence of silicate weathering upon respiration from roots and soil organisms. With these modifications, Caldeira and Kasting have extended the tenure of the biosphere to ten times that predicted by Lovelock and Whitfield: we can breathe easy for about a billion more years. But finally, temperatures will soar to the upper limits tolerable to the most primitive microbes, and in another billion years again the Earth will undergo a final sterilization with the

loss of the hydrogen from photolysed water to space.

Is an event so distant in time any cause for concern today? We certainly have more immediate problems during the coming century in trying to prevent our own selves from ravaging the biosphere. Nevertheless, Caldeira and Kasting's work raises issues that are hauntingly familiar: the role of CO<sub>2</sub> as a greenhouse gas, a possible need for planetary engineering in maintaining the Earth's habitability (or rather its comfort for humans), and fundamental questions about the ultimate persistence of life and the pace of evolution.

None of the models, of course, allows for possible evolutionary innovations that might buy time for life as we know it (or do not know it). The current plenitude of life's biochemical cycling mechanisms — within cells, organisms and ecosystems — is impressive, and has probably evolved in response to times of scarcity<sup>4</sup>. For example, C4 plants (such as maize) concentrate CO<sub>2</sub> in special bundle sheath cells, away from the sites of wasteful photorespiration; this innovation has evolved in the past 100 million years<sup>5</sup>, probably in response to falling levels of atmospheric CO<sub>2</sub>.

Future responses to scarcity could be technologically driven. In experiments for advanced life support, such as NASA's programme in closed ecological life-support systems, Biosphere 2 in Arizona, and the Russian Bios 3, carbon is assiduously cycled and CO<sub>2</sub> is managed at levels well above that of the Earth. Such research is intended to give us the rudiments of understanding that could someday lead to the propagation of mini-biospheres across the Galaxy. The future ancestral Earth — dry, burnt, and dead — would then be fondly recalled in Heinleinian space songs as that mythic place of "cool green hills".

Our descendants or descendent species would not have to run from the devolution, however — they could fight. Shades in space or mirrors on the Earth that keep out a small fraction of the elevated future solar flux would be an option. Carbonates could be heated to release CO<sub>2</sub>, a technique proposed<sup>6</sup> for the industrial stewardship of the carbon cycle on a 'terraformed' Mars<sup>7</sup> to make up for the lack of the CO<sub>2</sub>-generating plate tectonics that we have on Earth.

More importantly, long before the

solar-induced devolution of the biosphere comes to pass, the ability of *Homo sapiens* to persist will be tested not only by our own earthly transgressions, but by impacts from asteroids and comets. An impact that could destroy civilization and wipe out a quarter of the world's population is estimated to occur about every 500,000 years<sup>8</sup>.

Calculations of the life span of the biosphere bring to focus one additional matter of great philosophical import. If devolution-computing life forms evolved with just a billion years to spare, following a 3.5 billion year genealogy, we are presented with an intriguingly close temporal match between two essentially decoupled, developing systems. The time it takes for life and then technological intelligence to evolve is nearly identical, in the only case we know, to the window of opportunity afforded by an atmospheric resource, which is yoked, in part, to the luminosity history of a main-sequence star. That this window was open just long enough for our arrival will add impetus to the arguments used by proponents of the anthropic cosmological principle<sup>9</sup>, who find the Universe uncannily fine-tuned.

Overall, Caldeira and Kasting's study adds to our understanding of what Lovelock calls geophysiology and NASA calls Earth-system science: the Earth as an interconnected system of life, atmosphere, hydrosphere and lithosphere, a system whose organizational properties we are still discovering. While our concerns over habitability are often for the shorter term, in exploring the interplay of life and climate over geological time-scales, we forge an appreciation for where knowledge must go. **Tyler Volk**

*Tyler Volk is in the Earth Systems Group, Department of Applied Science, 26 Stuyvesant Street, New York University, New York 10003, USA.*

1. Caldeira, K. & Kasting, J. F. *Nature* **360**, 721–723 (1992).
2. Lovelock, J. E. & Whitfield, M. *Nature* **296**, 561–563 (1982).
3. Schwartzman, D. W. & Volk, T. *Palaeogeog. Palaeoclim. Palaeoecol.* **90**, 357–371 (1991).
4. Barlow, C. & Volk, T. *BioSystems* **23**, 371–384 (1990).
5. Ehleringer, J. R. et al. *Trends Ecol. Evol.* **6**, 95–99 (1991).
6. Fogg, M. J. *Br. Interplanet. Soc.* **45**, 315–329 (1992).
7. McKay, C. P. et al. *Nature* **352**, 489–496 (1992).
8. Morrison, D. *The Spaceguard Survey: Report of NASA International Near-Earth Object Detection Workshop* (January 25, 1992).
9. Carter, B. *Phil. Trans. R. Soc.* **A310**, 347–363 (1983).